Entanglement and coherence in distributed quantum networks

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Gdańsk



Outline

1 Quantum state merging and assisted entanglement distillation

2 Multipartite quantum state conversion

3 Assisted coherence distillation and incoherent state merging

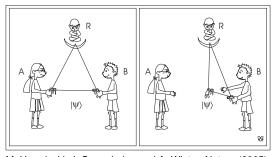
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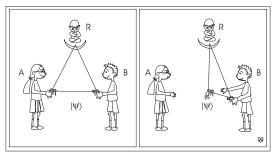
Quantum state merging



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Setting: Alice, Bob, and a referee share many copies of a pure state $|\psi\rangle^{RAB}$

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- Setting: Alice, Bob, and a referee share many copies of a pure state $|\psi\rangle^{RAB}$
- Aim of quantum state merging: send Alice's system to Bob while preserving the total state, i.e., the final state $|\psi\rangle^{RBB'}$ is the same as $|\psi\rangle^{RAB}$ up to relabeling A and B'

For achieving quantum state merging, Alice and Bob have access to shared singlets and a classical channel

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- The minimal number of singlets, asymptotically needed per copy of the state $|\psi\rangle^{RAB}$, is given by the conditional entropy:

$$S(A|B) = S(\rho^{AB}) - S(\rho^{B}) \tag{1}$$

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- Quantum conditional entropy can be <u>positive</u> or <u>negative</u>
- $S(A|B) \ge 0$: merging is possible with singlets at rate S(A|B)
- S(A|B) < 0: merging is possible without singlets, and Alice and Bob can obtain additional singlets at rate -S(A|B)

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Assisted entanglement distillation¹

■ Setting: Alice, Bob, and Charlie share many copies of a pure state $|\psi\rangle^{ABC}$

¹ D. P. DiVincenzo, C. A. Fuchs, H. Mabuchi, J. A. Smolin, A. Thapliyal, A. Uhlmann, Lecture Notes in Computer Science 1999; J. A. Smolin, F. Verstraete, A. Winter, PRA 2005

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- Setting: Alice, Bob, and Charlie share many copies of a pure state $|\psi\rangle^{ABC}$
- Aim of the process: asymptotic distillation of singlets between Alice and Bob by applying joint LOCC operations between all three parties

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Assisted entanglement distillation¹

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- Aim of the process: asymptotic distillation of singlets between Alice and Bob by applying joint LOCC operations between all three parties
- <u>Solution</u>: given a pure state $|\psi\rangle^{ABC}$, the optimal entanglement distillation rate between Alice and Bob with assistance of Charlie is equal to the regularized entanglement of assistance

$$E_a^{\infty}\left(\rho^{AB}\right) = \min\left\{S(\rho^A), S(\rho^B)\right\} \tag{2}$$

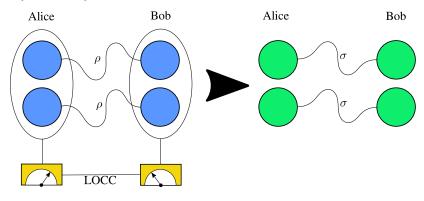
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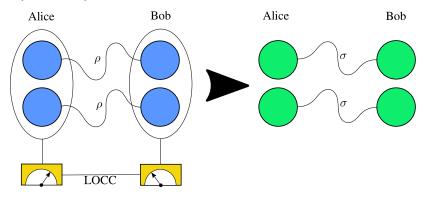
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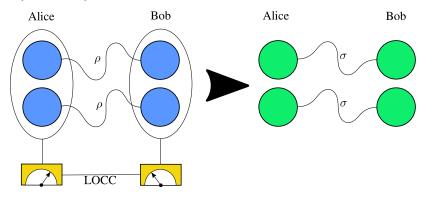
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- **Setting**: *N* parties share many copies of a multipartite state ρ
- Aim of multipartite state conversion: creation of another state σ via N-partite LOCC operations
- For N=2: entanglement distillation if $\sigma=|\Psi^+\rangle\langle\Psi^+|$, entanglement dilution if $\rho=|\Psi^+\rangle\langle\Psi^+|$

Conversion rate:

$$R(\rho \to \sigma) = \sup \left\{ r : \lim_{n \to \infty} \left(\inf_{\Lambda} \left\| \Lambda \left(\rho^{\otimes n} \right) - \sigma^{\otimes \lfloor rn \rfloor} \right\|_{1} \right) = 0 \right\} \quad (3)$$

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■ N = 2: $R(\rho \to |\Psi^+\rangle\langle\Psi^+|)$ is <u>distillable entanglement</u>¹ of ρ , and $R(|\Psi^+\rangle\langle\Psi^+|\to\sigma)^{-1}$ is <u>entanglement cost</u> of σ

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$$R(\psi^{AB} \to \phi^{AB}) = \frac{S(\psi^A)}{S(\phi^A)}$$
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■ Entanglement theory is <u>reversible</u> for bipartite pure states:

$$R(\psi^{AB} \to \phi^{AB}) = R(\phi^{AB} \to \psi^{AB})^{-1} \tag{5}$$

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■ Surprisingly little was known for N > 2

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- Surprisingly little was known for N > 2
- Bounds for N = 3:

$$R(\psi^{ABC} \to \phi^{ABC}) \le \min \left\{ \frac{S(\psi^A)}{S(\phi^A)}, \frac{S(\psi^B)}{S(\phi^B)}, \frac{S(\psi^C)}{S(\phi^C)} \right\}$$

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$$R(\psi^{ABC} \to \phi^{ABC}) \ge \min \left\{ \frac{S(\psi^A)}{S(\phi^B) + S(\phi^C)}, \frac{S(\psi^B)}{S(\phi^B)}, \frac{S(\psi^C)}{S(\phi^C)} \right\}$$
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■ The bound can be further improved by interchanging the parties *A*, *B*, *C*

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■ The bounds coincide whenever min $\left\{\frac{S(\psi^A)}{S(\phi^B)+S(\phi^C)}, \frac{S(\psi^B)}{S(\phi^B)}, \frac{S(\psi^C)}{S(\phi^C)}\right\}$ is equal to $\frac{S(\psi^B)}{S(\phi^B)}$ or $\frac{S(\psi^C)}{S(\phi^C)}$

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- In these cases we get the <u>exact rate</u> for tripartite quantum state conversion:

$$R(\psi^{ABC} \to \phi^{ABC}) = \min \left\{ \frac{S(\psi^B)}{S(\phi^B)}, \frac{S(\psi^C)}{S(\phi^C)} \right\}$$
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- Transformations among these states are in general not reversible: resource theory of entanglement is not reversible for N > 2

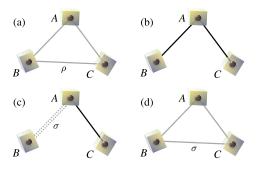
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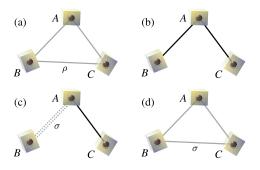
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- Transformations among these states are in general not reversible: resource theory of entanglement is not reversible for N > 2
- \blacksquare Results can be extended to N > 3

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Sketch of the proof:

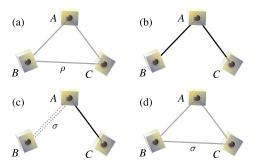
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Sketch of the proof:

 The parties apply <u>quantum state merging and assisted</u> <u>entanglement distillation</u> to distill singlets between AB and AC

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Sketch of the proof:

- The parties apply <u>quantum state merging and assisted</u> <u>entanglement distillation</u> to distill singlets between AB and AC
- Alice and Bob use their singlets to create the final state σ , remaining singlets between Alice and Charlie are used to teleport parts of σ to Charlie

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Incoherent states and measurements¹

<u>Incoherent states</u>: states which are diagonal in a preferred basis:

$$\sigma = \sum_{i} p_{i} |i\rangle \langle i| \tag{9}$$

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Incoherent states: states which are diagonal in a preferred basis:

$$\sigma = \sum_{i} p_{i} |i\rangle\langle i| \tag{9}$$

Incoherent measurements: quantum measurements which do not create coherence

$$\Lambda[\rho] = \sum_{i} K_{i} \rho K_{i} \tag{10}$$

with incoherent Kraus operators K_i , i.e., $K_i | m \rangle \sim | n \rangle$

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■ Maximally incoherent operations (MIO)¹: most general set, contains all operations which cannot create coherence: $\Lambda[\rho_i] \in \mathcal{I}$, where \mathcal{I} is the set of all incoherent states.

¹ J. Åberg, arXiv 2006

²A. Winter and D. Yang, PRL 2016; B. Yadin, J. Ma. D. Girolami, M. Gu, V. Vedral, PRX 2016

³G. Gour and R. W. Spekkens, NJP 2008

⁴B. Regula, M. Piani, M. Cianciaruso, T. R. Bromley, A. S., G. Adesso, arXiv 2017

⁵ M. Ringbauer, T. R. Bromley, M. Cianciaruso, S. Lau, G. Adesso, A. G. White, A. Fedrizzi, M. Piani, arXiv 2017

⁶A. S., G. Adesso, and M. B. Plenio, RMP 2017

- Maximally incoherent operations (MIO)¹: most general set, contains all operations which cannot create coherence: $\Lambda[\rho_i] \in I$, where I is the set of all incoherent states.
- Strictly incoherent operations (SIO)²: Incoherent operations for which also K_i^{\dagger} are incoherent. Correspond to quantum operations which do not use coherence.

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- Translationally invariant operations $(TIO)^3$: Quantum operations which commute with time translations, i.e., $e^{-iHt}\Lambda[\rho]e^{iHt} = \Lambda[e^{-iHt}\rho e^{iHt}]$ for a given Hamiltonian H.

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- Theory of multilevel coherence: coherence between N > 2 levels of a quantum system⁴⁵

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■ <u>Distillable coherence</u>¹: maximal rate for extracting the state $|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$ via incoherent operations

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- <u>Distillable coherence</u>¹: maximal rate for extracting the state $|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$ via incoherent operations
- $C_d(\rho) = S(\Delta[\rho]) S(\rho)$ with the dephasing operation Δ

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- The quantities differ for different frameworks of coherence
- Single-shot coherence distillation² and dilution³ has also been considered

¹ A. Winter and D. Yang, PRL 2016

²B. Regula, K. Fang, X. Wang, G. Adesso, arXiv 2017

³Q. Zhao, Y. Liu, X. Yuan, E. Chitambar, X. Ma, arXiv 2017

Incoherent operations on a single qubit admit a decomposition into 5 Kraus operators:

$$\left(\begin{array}{ccc} a_1 & b_1 \\ 0 & 0 \end{array}\right), \left(\begin{array}{ccc} 0 & 0 \\ a_2 & b_2 \end{array}\right), \left(\begin{array}{ccc} a_3 & 0 \\ 0 & b_3 \end{array}\right), \left(\begin{array}{ccc} 0 & b_4 \\ a_4 & 0 \end{array}\right), \left(\begin{array}{ccc} a_5 & 0 \\ 0 & 0 \end{array}\right)$$

¹ A. S., S. Rana, P. Boes, J. Eisert, PRL 2017

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a a_i are real, b_i are complex, $\sum_i a_i^2 = \sum_j |b_j|^2 = 1$, and $a_1b_1 + a_2b_2 = 0$

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- **a** a_i are real, b_i are complex, $\sum_i a_i^2 = \sum_j |b_j|^2 = 1$, and $a_1b_1 + a_2b_2 = 0$
- This characterization allows for complete solution of the single-qubit state conversion problem

¹ A. S., S. Rana, P. Boes, J. Eisert, PRL 2017

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- This characterization allows for complete solution of the single-qubit state conversion problem
- Open question: it is not known if 5 Kraus operators are indeed required, or if the number can be reduced to 4

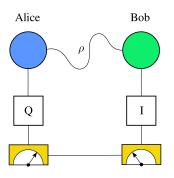
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Entanglement and coherence in distributed scenarios¹²³⁴

<u>LQICC</u>: Local quantum-incoherent operations and classical communication



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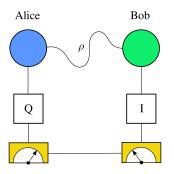
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LQICC operations preserve the set of quantum-incoherent states:

$$\rho_{\text{qi}}^{AB} = \sum_{i} p_{i} \sigma_{i}^{A} \otimes |i\rangle \langle i|^{B}$$

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■ <u>Setting</u>: Alice and Bob share many copies a bipartite state ρ^{AB} and can perform bipartite LQICC operations¹²

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- Setting: Alice and Bob share many copies a bipartite state ρ^{AB} and can perform bipartite LQICC operations¹²
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- For pure states $|\psi\rangle^{AB}$ optimal coherence distillation rate with assistance is $S(\Delta[\rho^B])$
- Optimal distillation rate without assistance: $S(\Delta[\rho^B]) S(\rho^B)$
- Confirmed in two recent experiments³⁴

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Standard quantum state merging: shared entanglement is a resource while local coherence is available at no cost

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- Optimal entanglement-coherence pairs (E, C): pairs of entanglement and coherence rates for which merging is possible, but neither E nor C can be reduced
- Main problem: determine all optimal pairs (E, C) for a given quantum state

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Theorem

Given a tripartite quantum state ρ^{RAB} , any achievable pair (E, C) fulfills the following inequality:

$$E + C \ge S(\Delta^{AB}[\rho^{RAB}]) - S(\Delta^{B}[\rho^{RAB}]).$$
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S is the von Neumann entropy and Δ^X denotes full decoherence of a (possibly multipartite) subsystem X:

$$\Delta^{X}[\rho] = \sum_{i} |i\rangle \langle i|^{X} \rho |i\rangle \langle i|^{X}.$$
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Since the right-hand side of Eq. (11) is nonnegative, the sum E+C is also nonnegative: no merging procedure can gain coherence and entanglement at the same time

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For pure states $|\psi\rangle^{RAB}$ we have

$$E \ge E_{\min} = S(\rho^{AB}) - S(\rho^{B}) \tag{13}$$

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The bound in Eq. (14) is achievable for all pure states

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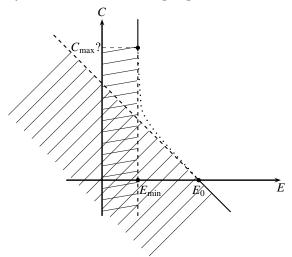
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Theorem

Any pure state $|\psi\rangle^{RAB}$ can be merged without local coherence by using singlets at rate

$$E_0 = S(\Delta^{AB}[\rho^{AB}]) - S(\Delta^B[\rho^B]). \tag{15}$$

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$$E_{\min} = S(\rho^{AB}) - S(\rho^{B}) \tag{16}$$

$$E_0 = S(\Delta^{AB}[\rho^{AB}]) - S(\Delta^B[\rho^B])$$
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■ Conjecture: it is possible to save a large amount of local coherence by using little extra entanglement, i.e., for some states the pairs $(E, C \gg 0)$ and $(E' = E + \varepsilon, C' \ll C)$ are both optimal for small ε

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- Possible candidate:

$$\rho = \frac{1}{d_B} \sum_{i=0}^{d_B-1} |i\rangle \langle i|^R \otimes |\phi_i\rangle \langle \phi_i|^A \otimes |\psi_i\rangle \langle \psi_i|^B, \qquad (18)$$

where $|\psi_i\rangle$ are mutually orthogonal maximally coherent states of arbitrary dimension d_B , and $|\phi_i\rangle$ are single-qubit states

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- We found lower and upper bounds for <u>multipartite state</u> <u>conversion rates</u>
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- Assisted distillation of coherence has also been performed in two recent experiments
- We introduced the task of incoherent quantum state merging, in which both entanglement and local coherence are considered as a resource
- Our results imply an incoherent version of Schumacher compression: $S(\Delta[\rho])$ is the optimal compression rate if the decompression is performed via incoherent operations only

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