

Multi-partite entanglement can speed up quantum key distribution in networks

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Outline

- Entanglement-based quantum key distribution (QKD)

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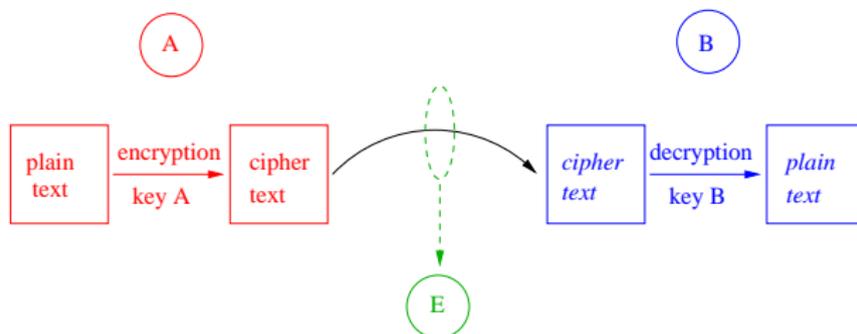
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*M. Epping, H. Kampermann, C. Macchiavello, and DB, New J. Phys. **19**, 093012 (2017)*

Quantum key distribution (QKD)



Vernam cipher \equiv “one-time pad” (1917):

Encoding with secret random key (only known to Alice and Bob, not to Eve). Proven to be secure.

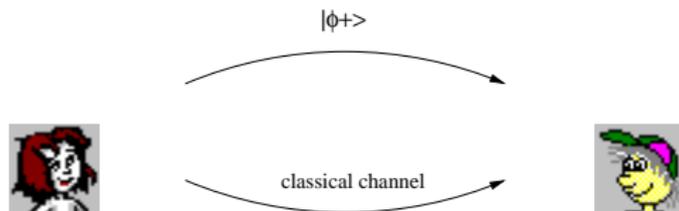
How to establish secret random key?

\hookrightarrow quantum cryptography \equiv quantum key distribution (QKD)

Entanglement-based QKD (between two parties)

A. Ekert, *Phys. Rev. Lett.* **67**, 661 (1991)

Aim: secret random key for Alice and Bob



1) A sends half of a Bell state to Bob: $|\phi^+\rangle_{AB} = \frac{1}{\sqrt{2}}(|00\rangle_{AB} + |11\rangle_{AB})$
A and B measure, use 2 bases randomly: \uparrow or \searrow

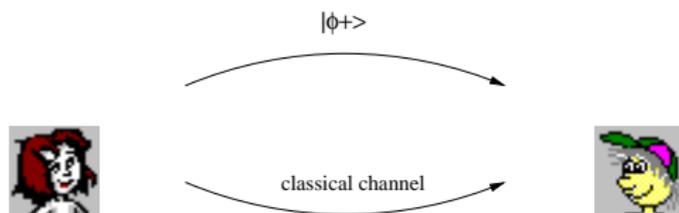
2) A and B exchange class. info about basis,
keep matching cases: $1 \ r \ 0 \ 0 \ 1 \ r \ 0 \ r$

↪ Alice and Bob have established secret random key!

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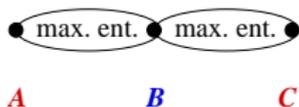
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Security: monogamy of entanglement

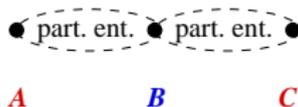


Monogamy of entanglement

V. Coffman, J. Kundu, and W. K. Wootters, *Phys. Rev. A* **61**, 052306 (2000)



Impossible!



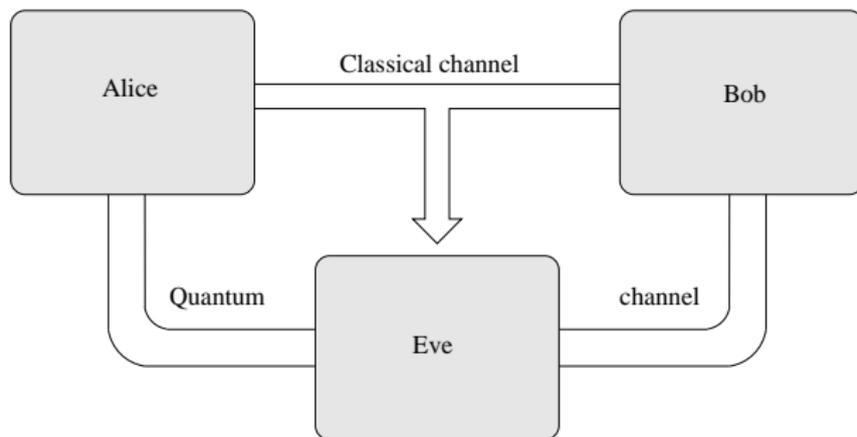
Possible

$$E(B|A) + E(B|C) \leq E(B|AC)$$

QKD in reality: noisy entangled state, $\rho = p|\phi^+\rangle\langle\phi^+| + (1-p)\frac{1}{4}\mathbf{1}$,
assume Eve to have purifying state (is partially correlated with A/B)

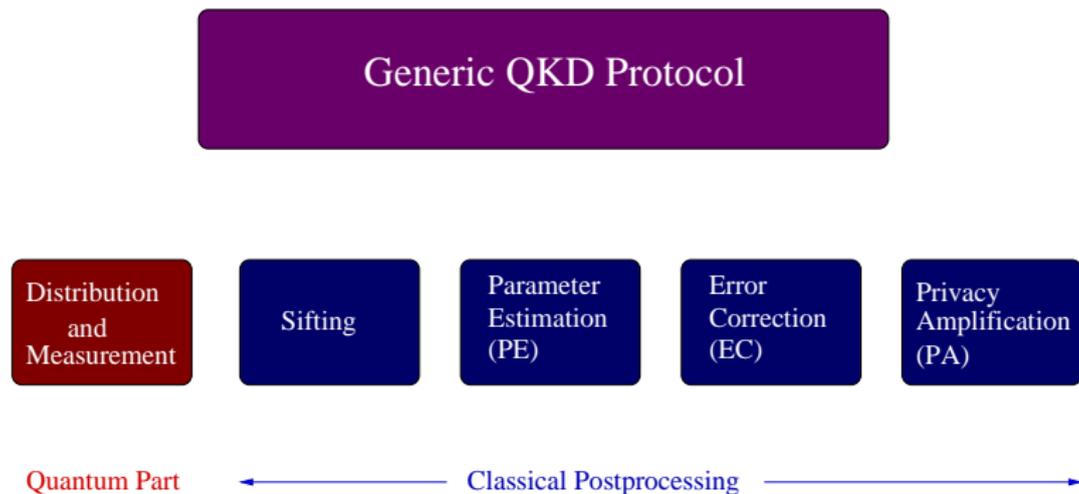
↪ security analysis

Quantum Key Distribution (QKD)

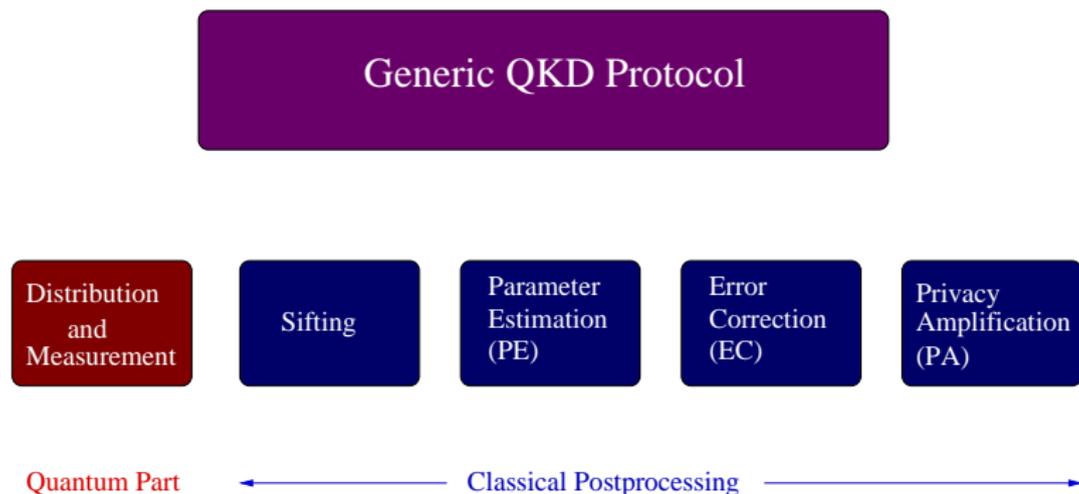


- Scenario: Alice und Bob have quantum channel (controlled by Eve) and classical channel (authenticated)
- Secure communication \Leftrightarrow Creation of a secret random key pair between Alice and Bob
- No restrictions on Eve

QKD: General description of a QKD protocol



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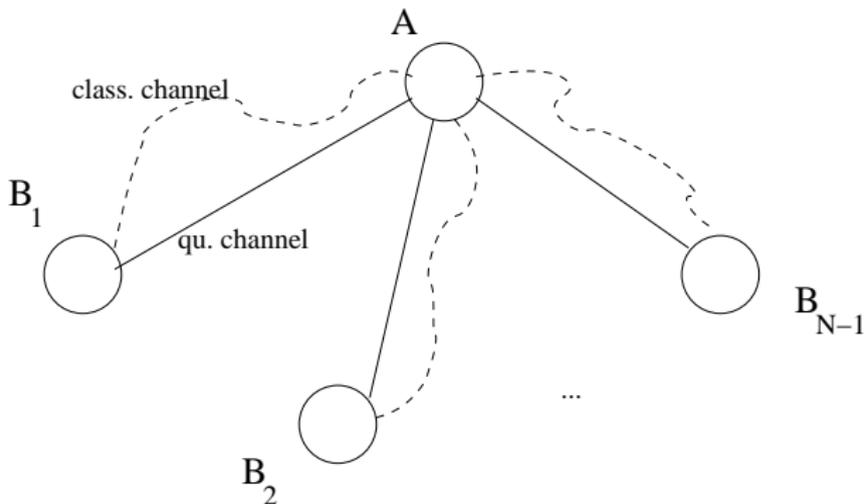
Equivalence of prepare+measure QKD with entanglement-based QKD

↔ In the following: use entanglement-based scheme

Generalisation of QKD to more than two parties

M. Epping, H. Kampermann, C. Macchiavello, and DB, New J. Phys. 19, 093012 (2017)

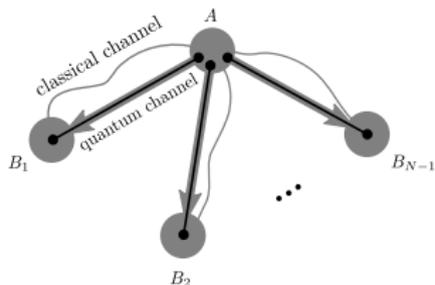
Aim: establish joint secret random key between N parties,
i.e. “conference key”



Establishing a conference key: Two possibilities

*M. Epping, H. Kampermann, C. Macchiavello, and DB, New J. Phys. **19**, 093012 (2017)*

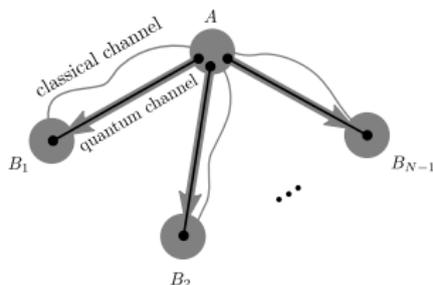
Using bipartite entanglement (2QKD):



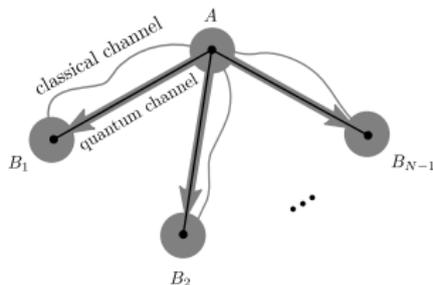
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Using bipartite entanglement (2QKD):



... or using multipartite entanglement (NQKD):



Multipartite entanglement

Multipartite entanglement

Multipartite entanglement of composite (pure) states of N parties:

$$|\psi\rangle = |a\rangle_{1,\dots,k} \otimes |b\rangle_{k+1,\dots,N} \iff \text{separable across bipartite split}$$

$$|\psi\rangle \neq |a\rangle_{1,\dots,k} \otimes |b\rangle_{k+1,\dots,N} \iff \text{multipartite entangled}$$

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Example (separable): $|\psi\rangle = |0\rangle|0\rangle\dots|0\rangle$

Example (entangled): **GHZ states of N qubits**

$$|\psi_j^\pm\rangle = \frac{1}{\sqrt{2}}(|0\rangle|j\rangle \pm |1\rangle|\bar{j}\rangle)$$

where j takes values $0, \dots, 2^{N-1} - 1$ in binary notation;
 \bar{j} is negation of j , e.g. if $j = 010$ then $\bar{j} = 101$

Multipartite entanglement for QKD

Which types of multipartite entanglement can be used for QKD?

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Theorem (Perfect resource state for multipartite QKD)

For N qubits, with $N \geq 3$, the state

$|\phi_{corr}\rangle = a_{0,\dots,0}|0, \dots, 0\rangle + a_{1,\dots,1}|1, \dots, 1\rangle$ with $|a_{0,\dots,0}|^2 + |a_{1,\dots,1}|^2 = 1$

leads to perfect classical correlations between any number of parties, if and only if each of them measures in the z -basis.

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Proof: “ \Leftarrow ” clear;

“ \Rightarrow ”: observable \mathcal{M}_{ij} of two parties i and j :

$$\mathcal{M}_{ij} = (\vec{M}_i \cdot \vec{\sigma}) \otimes (\vec{M}_j \cdot \vec{\sigma}) = \sum_{\alpha, \beta \in \{x, y, z\}} M_i^\alpha M_j^\beta \sigma_i^\alpha \otimes \sigma_j^\beta,$$

$$\langle \phi_{corr} | \sigma_i^\alpha \otimes \sigma_j^\beta | \phi_{corr} \rangle = 0 \quad \text{unless } \alpha = \beta = z,$$

also $\langle \phi_{corr} | \sigma_i^\alpha \otimes \sigma_j^\beta | \phi_{corr} \rangle = 2[p_i^\alpha(+)p_j^\beta(+) + p_i^\alpha(-)p_j^\beta(-)] - 1$,

thus $p_i^\alpha(+)p_j^\beta(+) + p_i^\alpha(-)p_j^\beta(-) \neq 1$, unless $\alpha = \beta = z$.

Multipartite QKD protocol

If one requires perfect correlations and uniformity of key, the *only* possible resource state is $|GHZ\rangle = \frac{1}{\sqrt{2}}(|0, \dots, 0\rangle + |1, \dots, 1\rangle)$.

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- 1) *State preparation:* Parties A and B_i , $i = 1, 2, \dots, N - 1$ share $|GHZ\rangle = \frac{1}{\sqrt{2}} \left(|0\rangle^{\otimes N} + |1\rangle^{\otimes N} \right)$.

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- 4) *Classical post-processing*: As in the bipartite protocol, error correction and privacy amplification is performed.

Secret key rate for NQKD

Security analysis:

- Analogous to bipartite case, with modifications in worst-case error correction and depolarisation

R. Renner, N. Gisin, and B. Kraus, Phys. Rev. A 72, 012332 (2005)

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- Figure of merit: **secret fraction**,
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$$r_\infty = \sup_{U \leftarrow K} \inf_{\sigma_{A\{B_i\}} \in \Gamma} [S(U|E) - \max_{i \in \{1, \dots, N-1\}} H(U|K_i)],$$

with $U \leftarrow K$: bitwise preprocessing channel on A 's raw key bit K ,

$S(U|E)$: conditional von-Neumann entropy of (class.) key variable and E ,

$H(U|K_i)$: conditional Shannon entropy of U and B_i 's guess of it,

Γ : set of all density matrices $\sigma_{A\{B_i\}}$ of A and B_i consistent with parameter estimation

Secret key rate: $R = r_\infty R_{\text{rep}}$ with repetition rate R_{rep}

Secret key rate for NQKD

Introduce (extended) depolarisation procedure, \leftrightarrow GHZ-diagonal state
 \leftrightarrow calculate **secret fraction** r_∞ :

Secret key rate for NQKD

Introduce (extended) depolarisation procedure, \hookrightarrow GHZ-diagonal state
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$$\begin{aligned} r_\infty = & \left(1 - \frac{Q_Z}{2} - Q_X\right) \log_2 \left(1 - \frac{Q_Z}{2} - Q_X\right) \\ & + \left(Q_X - \frac{Q_Z}{2}\right) \log_2 \left(Q_X - \frac{Q_Z}{2}\right) \\ & + (1 - Q_Z)(1 - \log_2(1 - Q_Z)) - h\left(\max_{1 \leq i \leq N-1} Q_{AB_i}\right) \end{aligned}$$

with Q_Z : probability that at least one B_i obtains different result than A in z -measurement,
with Q_X : probability that at least one B_i obtains in x -measurement a result that is
incompatible with noiseless state,

binary entropy: $h(p) = -p \log_2 p - (1 - p) \log_2 (1 - p)$,

Q_{AB_i} : probability that z -measurements of A and B_i disagree.

Example for explicit key rates

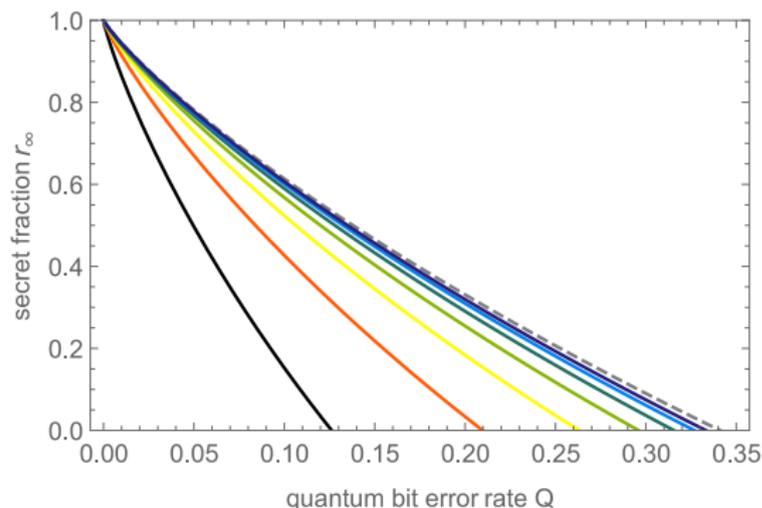
Noise model: mixture of GHZ-state and white noise (then $Q = Q_z$)

$$r_{\infty}(Q, N) = 1 + h(Q) - h\left(Q \frac{2^N - 1}{2^N - 2}\right) - h\left(Q \frac{2^{N-1}}{2^N - 2}\right) \\ + \left(\log_2(2^{N-1} - 1) - \frac{2^N - 1}{2^N - 2} \log_2(2^N - 1)\right) Q,$$

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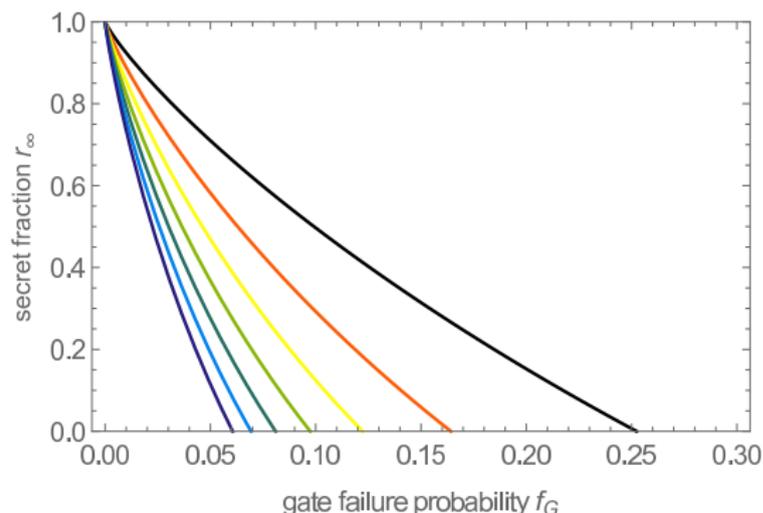
Key rates for $N = 2, 3, \dots, 8$,
from left to right.

Secret key rate as function of gate failure probability

Consider imperfect state preparation (depolarising noise): experimental creation of GHZ-state is more demanding with higher N !

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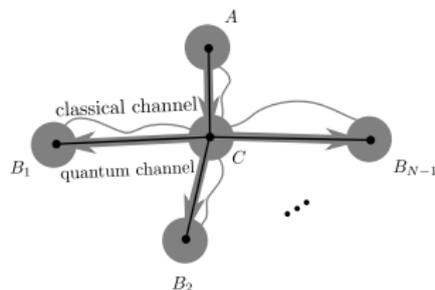
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Key rates for $N = 2, 3, \dots, 8$,
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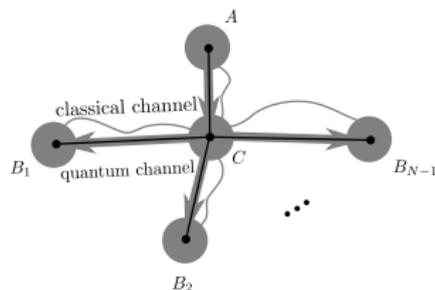
Advantage of NQKD in quantum networks

Consider quantum networks with routers (can produce and entangle qubits), fixed channel capacity:

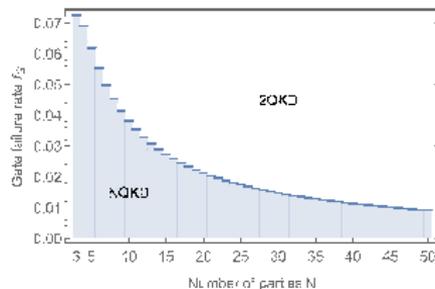


Advantage of NQKD in quantum networks

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For small gate failure probability: NQKD is better than 2QKD!



Connection to quantum network coding

Processing of data at intermediate network nodes can improve throughput and increase robustness of quantum network with bottleneck.

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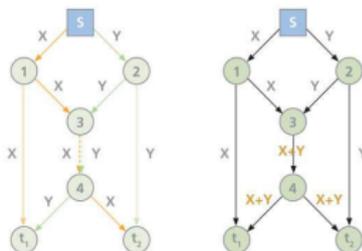
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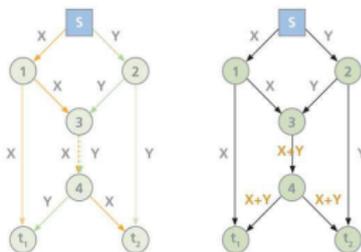


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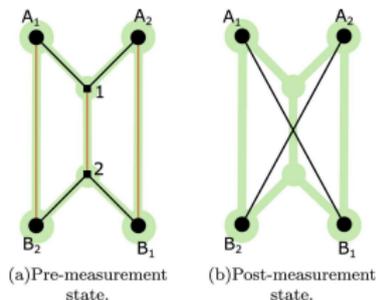
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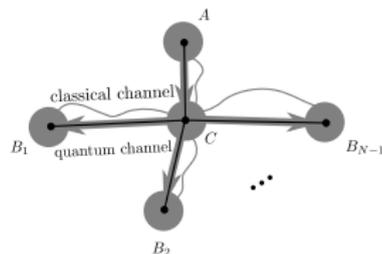
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M. Epping, H. Kampermann, and DB, New J. Phys. 18, 103052 (2016)

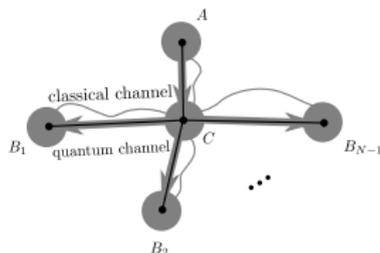
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Distribution of GHZ-state in above network, with quantum operations at node C (router), and fixed channel capacities for all links:



Connection to quantum network coding

Distribution of GHZ-state in above network, with quantum operations at node C (router), and fixed channel capacities for all links:



- A produces Bell state and sends only one qubit C to router:
$$|\text{---}\rangle_{CA} = \frac{1}{\sqrt{2}}(|0+\rangle + |1-\rangle)_{CA}$$
- C produces $(N - 1)$ qubits and entangles them with C via C_z gates:
$$|\psi_{\text{total}}\rangle = \frac{1}{\sqrt{2}}(|+\rangle_C |GHZ'\rangle_{AB_i} + |-\rangle_C X_{B_1} |GHZ'\rangle_{AB_i})$$

where $|GHZ'\rangle$ is GHZ-state in X -basis.
- Router measures qubit C in X -basis and distributes qubits to B_i .
- Impossible to create $(N - 1)$ Bell pairs by sending single qubit from A to router; need $(N - 1)$ network uses.

M. Epping, H. Kampermann, and DB, *New J. Phys.* **18**, 103052 (2016)

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Quantum Information Theory in Düsseldorf

Institut für Theoretische Physik III, Heinrich-Heine-Universität Düsseldorf, Germany



from left to right: J. Bremer, J. M. Henning, D. Miller, H. Kampermann, T. Holz, G. Gianfelici, M. Zibull, DB, T. Backhausen, S. Datta, F. Bischof, T. Wagner, C. Liorni, C. Glowacki, F. Grasselli, C. Hoffmeister, B. Sanvee, L. Tendick, M. Battiato

