

The Role of Raychaudhuri equation in a Self Similar Collapse

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Raychaudhuri equation

Raychaudhuri equation for a timelike congruence having four velocity u^μ is given by,

[A. K. Raychaudhuri (1955), J. Ehlers (1993)]

$$\frac{d\theta}{d\lambda} = -\frac{1}{3}\theta^2 + \nabla_\mu a^\mu - \sigma_{\mu\nu}\sigma^{\mu\nu} + \omega_{\mu\nu}\omega^{\mu\nu} - R_{\mu\nu}u^\mu u^\nu$$

$$\theta = \nabla_\mu u^\mu$$

$$\sigma_{\mu\nu} = \nabla_{[\nu} u_{\mu]} - \frac{1}{3}h_{\mu\nu}\theta + a_{[\nu} u_{\mu]}$$

$$\omega_{\mu\nu} = \nabla_{[\nu} u_{\mu]} - a_{[\nu} u_{\mu]}$$

$$a_\mu = u^\nu \nabla_\nu u_\mu$$

$$h_{\mu\nu} = g_{\mu\nu} + u_\mu u_\nu$$

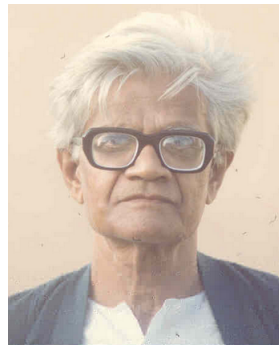
expansion scalar

shear tensor

rotation tensor

four acceleration

spatial metric



Raychaudhuri equation

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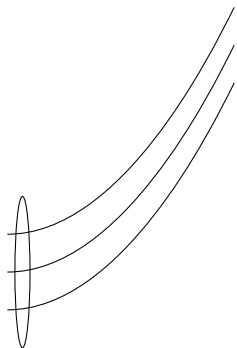


Figure : Area enclosing a congruence

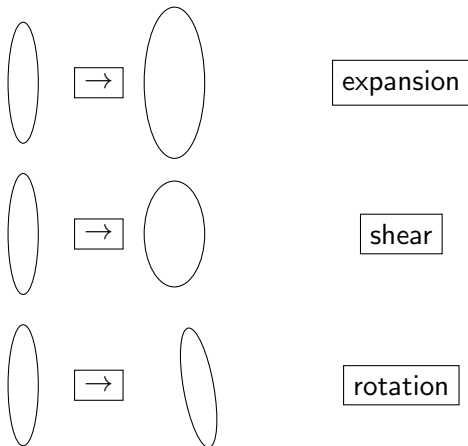


Figure : Expansion, Shear and Rotation

[S. Kar & S. Sengupta (2007)]

Metric and Self-Similarity

Metric and Self-Similarity

- The metric is assumed to be conformally flat -

$$ds^2 = \frac{1}{A(r, t)^2} \left[- dt^2 + dr^2 + r^2 d\Omega^2 \right],$$

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- The energy-momentum tensor is that of a massless scalar -

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- A spacetime is assumed to be self-similar such that,

[A. H. Cahill & M. E. Taub, (1971)],

[B.J. Carr & A.A. Coley, (1999)]

$$A(r, t) = B(z), \quad \phi(r, t) = \phi(z), \quad z = t/r$$

Analysis of the behavior

- For a comoving observer- $u^\alpha = A\delta_0^\alpha$, with the energy-momentum tensor of a massless scalar field, Raychaudhuri equation can be written as,

$$\frac{d\theta}{d\lambda} + \frac{1}{3}\theta^2 + \left(-3A'^2 + \frac{2AA'}{r} + AA'' \right) = A^2\dot{\phi}^2$$

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- Writing $A(r, t) = B(z)$, where $z = t/r$, for self similar case, the above condition becomes,

$$\left(\frac{d\phi}{dz}\right)^2 \geq \frac{3z^2}{B^2} \left(\frac{dB}{dz}\right)^2 - \frac{z^2}{B} \frac{d^2B}{dz^2}$$

Analysis of the behavior

- Using information from field equations, the condition for focusing can be written as,

$$\frac{1}{B} \frac{d^2 B}{dz^2} + P(z) \frac{1}{B^2} \left(\frac{dB}{dz} \right)^2 + Q(z) \frac{1}{B} \frac{dB}{dz} \begin{cases} \geq 0 & \text{for } 0 < z < 1 \\ \leq 0 & \text{for } 1 < z < 2 \\ \geq 0 & \text{for } z > 2 \end{cases}$$

$$P(z) = \frac{3z^3 - 3z^2 - 6z - 6}{4 + z^2 - z^3}, \quad Q(z) = \frac{4(2z - 1)}{(1 - z)(4 + z^2 - z^3)}$$

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- This condition is studied for high z values which is chosen in the limit $t \rightarrow \text{finite}$ and $r \rightarrow 0$ (Central Singularity).

- For high values of z , the condition for focusing can be approximated as,

$$\left(\frac{3}{z} + \frac{8}{z^3}\right) \frac{1}{B} \frac{dB}{dz} \geq 0$$

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- Thus, for focusing $\frac{dB}{dz}$ has to be non-negative.
- The situation is further analyzed by solving the condition -

$$\frac{1}{B} \frac{d^2B}{dz^2} + P(z) \frac{1}{B^2} \left(\frac{dB}{dz}\right)^2 + Q(z) \frac{1}{B} \frac{dB}{dz} \geq 0,$$

for the limiting case.

Analysis of the behavior

- Variation of $\frac{dB(z)}{dz}$ with a positive initial $\frac{dB}{dz}$:

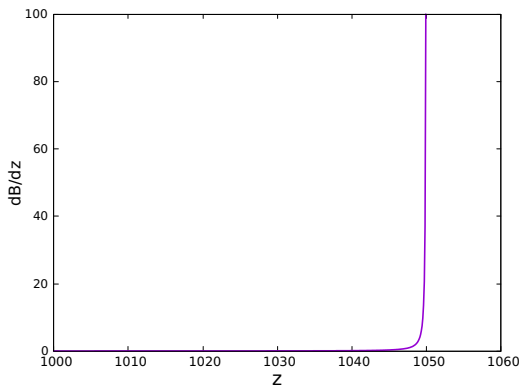


Figure : Plot of $\frac{dB}{dz}$ with z ($B_0 = 1$, $(\frac{dB}{dz})_0 = 0.01$)

Analysis of the behavior

- Variation of $B(z)$ with a positive initial $\frac{dB}{dz}$:

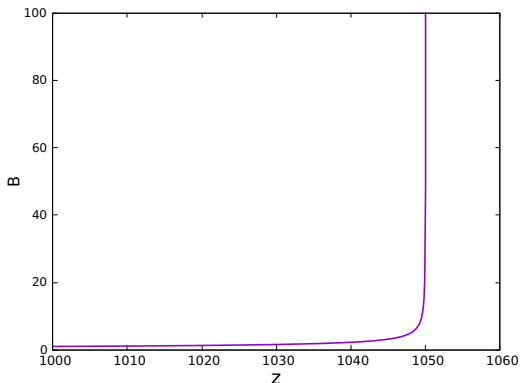


Figure : Plot of B with z ($B_0 = 1$, $(\frac{dB}{dz})_0 = 0.01$)

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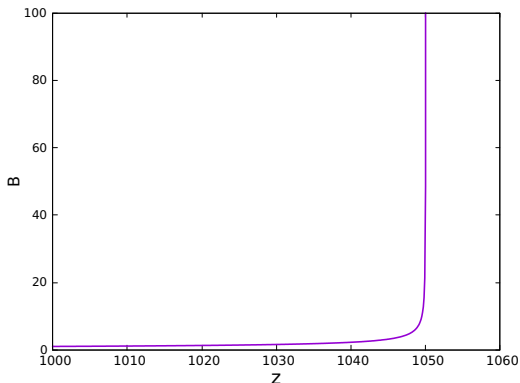


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- For this kind of initial condition the scalar field collapses and a zero proper volume singularity forms.

Analysis of the behavior

- Variation of $\frac{dB(z)}{dz}$ with a negative initial $\frac{dB}{dz}$:

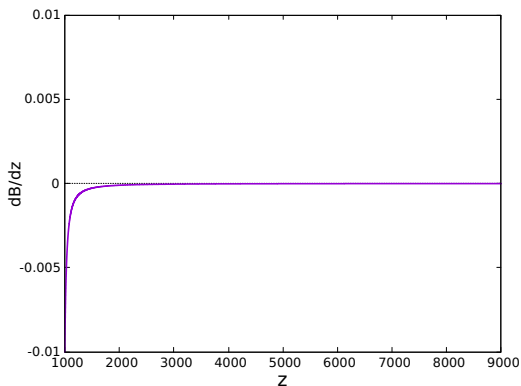


Figure : Plot of $\frac{dB}{dz}$ with z ($B_0 = 1$, $(\frac{dB}{dz})_0 = -0.01$)

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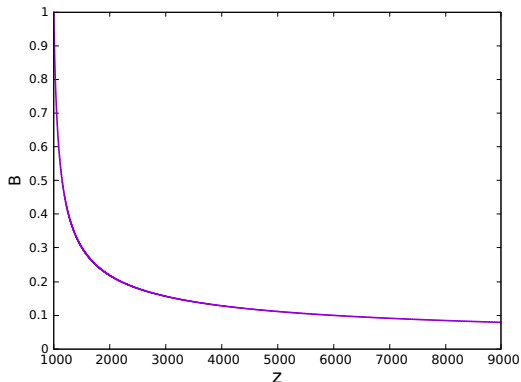


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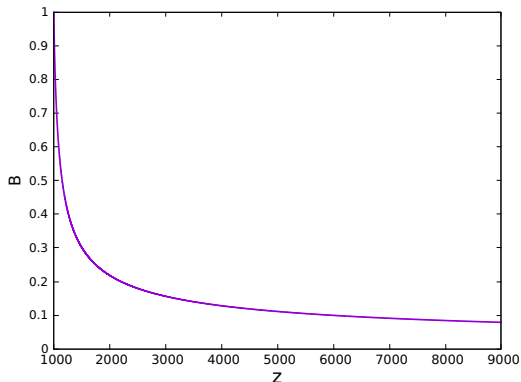


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- The solution is non-collapsing for this kind of initial condition.

Approximate analytic study

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- The condition which ensures focusing is -

$$\frac{1}{B} \frac{d^2 B}{dz^2} + P(z) \frac{1}{B^2} \left(\frac{dB}{dz} \right)^2 + Q(z) \frac{1}{B} \frac{dB}{dz} \geq 0$$

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- If u_{eq} and B_{eq} are the solutions for $\frac{du}{dz} = 2u^2$,

$$\frac{1}{u_{\text{eq}}} = C - 2z, \quad B_{\text{eq}} = \frac{D}{\sqrt{|C - 2z|}}, \quad R_{\text{eq}}(r, t) = \frac{r^{\frac{1}{2}} |Cr - 2t|^{\frac{1}{2}}}{D}.$$

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- If initially $u = u_{\text{eq}}$ and $u, u_{\text{eq}} > 0$, $u \geq u_{\text{eq}}$ always. Thus $u_{\text{eq}} \rightarrow \infty$ ensures $u \rightarrow \infty$. Similarly, $B \geq B_{\text{eq}}$ always and the solution is collapsing.

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- For second set of initial data $u, u_{\text{eq}} < 0$ initially. Thus $u \leq u_{\text{eq}}$ and $B \leq B_{\text{eq}}$ always.
- It should be noted that the condition for singularity formation has already been established using Raychaudhuri equation as,

$$\frac{dB}{dz} \geq 0$$

Approximate analytic study

- For the collapsing solution, few more investigations are performed using,

$$B = B_{\text{eq}} = \frac{B_0}{\sqrt{C - 2z}}$$

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$$B = B_{\text{eq}} = \frac{B_0}{\sqrt{C - 2z}}$$

- Singularity forms at $z = C/2$. Thus, innermost shells become singular before the outer ones.
- Using the condition for shell crossing singularity,

[P. Yodzis et al. (1973), B. Waugh & K. Lake (1986)]

$$\frac{\partial g_{\theta\theta}}{\partial r} = 0, \quad g_{\theta\theta} \neq 0, \quad r > 0$$

it is proved that this is not a shell crossing singularity, rather it is shell focusing.

- The study of apparent horizon ($g^{\mu\nu} R_{,\mu} R_{,\nu} = 0$, $R = r/A(r, t)$) shows that the singularity forms at a smaller z value than the apparent horizon which is the signature of a Naked singularity.

Approximate analytic study

- The study of apparent horizon ($g^{\mu\nu} R_{,\mu} R_{,\nu} = 0$, $R = r/A(r, t)$) shows that the singularity forms at a smaller z value than the apparent horizon which is the signature of a Naked singularity.
- The expression for the scalar field is given by,








$$\phi = \ln \left[\frac{F}{(C - 2z)^{E/2}} \right] \text{ for } C > 2z$$
$$\phi = \ln \left[F (2z - C)^{E/2} \right] \text{ for } C < 2z$$

where $E > 0$.

Conclusions and Future Work

- We have studied the evolution of conformally flat self similar spacetime minimally coupled to a massless scalar field using Raychaudhuri equation.
- We have found the dependence of the evolution on the different type of initial conditions which may give rise to collapsing or regular solution.
- Different properties and nature of the final outcome have been analyzed for the collapsing case.
- Using this approach useful information about the spacetime can be extracted without explicitly solving the field equations.
- Our immediate future goal is to consider the massive scalar field case.
- We want to analyze more general spacetimes using this approach to see its usefulness.

References

-  A. H. Cahill and M. E. Taub, *Comm. Math. Phys.* **21**, 1 (1971).
-  B.J. Carr and A.A. Coley, *Class. Quantum Grav.*, **16**, R31 (1999).
-  A. K. Raychaudhuri, *Phys. Rev.*, **98**, 1123 (1955).
-  J. Ehlers, *Gen. Relativ. Gravit.* **25**, 1225 (1993).
-  S. Kar and S. Sengupta, *Pramana-J Phys* **69**: 49 (2007).
-  P. Yodzis, H.-J. Seifert and H. Muller zum Hagen, *Commun. math. Phys.* **34**, 135 (1973).
-  B. Waugh and K. Lake, *Phys. Rev. D* **34**, 2978 (1986).

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Thank You