

The curious case of torsion

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What and why?

- 1 Torsion is set to zero in GR
- 2 Palatini variation (metric and connection independent) confirms torsion-free condition
- 3 vierbein-Einstein-Palatini (VEP)/Einstein-Cartan-Sciama-Kibble (ECSK) theory
- 4 Fermions, torsion and conformal transformation

- Gravity is described in terms of independent variables (Palatini type) tetrads e_μ^I and spin connection A_μ^{IJ} [Kibble (1961), Sciama (1964), Hehl and Datta (1971)]

$$g_{\mu\nu} = e_\mu^I e_\nu^J \eta_{IJ}, \Gamma_{\mu\nu}^\alpha = e_I^\alpha \partial_\mu e_\nu^I + A_{\mu J}^I e_\nu^J e_I^\alpha \quad (1)$$

- Torsion

$$C_{\alpha\mu}^\alpha = \Gamma_{\mu\nu}^\alpha - \Gamma_{\nu\mu}^\alpha \quad (2)$$

- Action

$$S = \frac{1}{2\kappa} \int |e| d^4x F_{\mu\nu}^{IJ} + S_M \quad (3)$$

- Equations

$$\delta e : F_{\alpha\mu}^{IJ} e_I^\alpha - \frac{1}{2} e_\mu^J F_{\alpha\beta}^{KL} e_K^\alpha e_L^\beta = \kappa \Theta_\mu^J \quad (4)$$

$$\delta A : A_\mu^{IJ} = \omega_\mu^{IJ}(e) + \Lambda_\mu^{IJ} \quad (5)$$

- Torsion

$$C_{\mu\nu}^\alpha = e_I^\alpha \Lambda_{[\mu}^{IJ} e_{\nu]J} \quad (6)$$

- Fermion (ψ) [Hehl and Datta (1971)]

Spinor covariant derivative:

$${}^{\psi}D_{\mu}\psi = \partial_{\mu}\psi - \frac{i}{4}A_{\mu}^{IJ}\sigma_{IJ}\psi. \quad (7)$$

Action:

$$\begin{aligned} S[e, A, \phi, \psi] = S[e, A] + \int |e| d^4x \frac{i}{2} & \left[\left(\bar{\psi} \gamma^K e_K^{\mu} {}^{\psi}D_{\mu}\psi \right. \right. \\ & \left. \left. - (\bar{\psi} \gamma^K e_K^{\mu} {}^{\psi}D_{\mu}\psi)^{\dagger} \right) \right] \end{aligned} \quad (8)$$

- On-shell spin connection

$$A_{\mu}^{IJ} = \omega_{\mu}^{IJ} + \frac{\kappa}{8} \bar{\psi} \{ \gamma_K, \sigma^{IJ} \} \psi e_{\mu}^K \quad (9)$$

- On-shell torsion

$${}^{OS}C^{\alpha}_{\mu\nu} = \frac{\kappa}{2} \epsilon^{IJKL} \bar{\psi} \gamma_L \gamma_5 \psi e_I^{\alpha} e_{J\mu} e_{K\nu}. \quad (10)$$

- Non-linear Dirac equation (NLD)

$$\gamma^K e_K^{\mu} \partial_{\mu} \psi + \frac{1}{2} C^{\alpha}_{\mu\alpha} e_K^{\mu} \gamma^K \psi - \frac{i}{4} A_{\mu}^{IJ} e^{\mu K} \gamma_K \sigma_{IJ} \psi = 0. \quad (11)$$

If C is identified with ${}^{OS}C$:

$$\gamma^K e_K^{\mu} \partial_{\mu} \psi - \frac{i}{4} \omega_{\mu}^{IJ} e^{\mu K} \gamma_K \sigma_{IJ} \psi - \frac{i\kappa}{64} \bar{\psi} \{ \gamma_K, \sigma_{IJ} \} \psi \{ \gamma^K, \sigma^{IJ} \} \psi = 0. \quad (12)$$

- With *a priori* torsion-free condition

$$\gamma^K e_K^{\mu} \partial_{\mu} \psi - \frac{i}{4} \omega_{\mu}^{IJ} e^{\mu K} \gamma_K \sigma_{IJ} \psi = 0. \quad (13)$$

- Einstein's equation

$$\widehat{R}_{\mu\nu} - \frac{1}{2}g_{\mu\nu}\widehat{R} = \kappa\widehat{T}_{\mu\nu}(\psi, \bar{\psi}) - \frac{3\kappa^2}{16}g_{\mu\nu}\bar{\psi}\gamma_I\gamma_5\psi\bar{\psi}\gamma^I\gamma_5\psi. \quad (14)$$

- Effective Lagrangian in torsion-free theory

$$\begin{aligned} \mathcal{L}_F = & \frac{i}{2} \left(\bar{\psi}\gamma^K e_K^\mu \partial_\mu \psi - \partial_\mu \bar{\psi}\gamma^K e_K^\mu \psi - \frac{i}{2} \epsilon_{IJKL} \omega_\mu^{IJ} e^{\mu K} \bar{\psi}\gamma^L \gamma_5 \psi \right. \\ & \left. - \frac{3i\kappa}{8} ((\bar{\psi}\gamma_5\psi)^2 - (\bar{\psi}\psi)^2) \right). \end{aligned} \quad (15)$$

Conformal transformation

- Torsion-free Dirac equation is invariant; NLD is not under conformal transformation
 $(g_{\mu\nu} \rightarrow \Omega^2 g_{\mu\nu}, \tilde{e}_\mu^I = \Omega e_\mu^I, \psi \rightarrow \Omega^{-\frac{3}{2}} \psi)$ [Chakrabarty and Lahiri (2018)]

$$\Omega^{-\frac{5}{2}} \left(\gamma^K e_K^\mu \partial_\mu \psi - \frac{i}{4} \omega_\mu^{IJ} e^{\mu K} \gamma_K \sigma_{IJ} \psi \right) - \frac{i\kappa}{64} \Omega^{-\frac{9}{2}} \bar{\psi} \{ \gamma_K, \sigma_{IJ} \} \psi \{ \gamma^K, \sigma^{IJ} \} \psi = 0. \quad (16)$$

- How should torsion (spin connection) behave under conformal transformation?
 - Off-shell:
 - Nieh-Yan theory [Nieh and Yan 1982] $A_\mu^{IJ} \rightarrow A_\mu^{IJ}$
 - Invariant torsion $A_\mu^{IJ} \rightarrow A_\mu^{IJ} + (e_\mu^I e^{J\alpha} - e_\mu^J e^{I\alpha}) \partial_\alpha \ln \Omega$
 - On-shell: Dynamically generated torsion \rightarrow NLD

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Conformal properties of torsion

- General off-shell transformation

$$A_{\mu}^{IJ} \rightarrow A_{\mu}^{IJ} + \xi \left(e_{\mu}^I e^{J\nu} - e_{\mu}^J e^{I\nu} \right) \partial_{\nu} \ln \Omega \quad (17)$$

$$\Lambda_{\mu}^{IJ} \rightarrow \Lambda_{\mu}^{IJ} - (1 - \xi) \left(e_{\mu}^I e^{J\nu} - e_{\mu}^J e^{I\nu} \right) \partial_{\nu} \ln \Omega, \quad (18)$$

$\xi = 0$: Nieh-Yan theory

$\xi = 1$: Invariant torsion

- General conformal transformation of torsion

$$C^{\alpha}_{\mu\nu} \rightarrow C^{\alpha}_{\mu\nu} + (\xi - 1) \delta^{\alpha}_{[\mu} \partial_{\nu]} \ln \Omega \quad (19)$$

- Dirac equation is invariant with the general transformations

- Gravity is different for different ξ !

$$\begin{aligned} R \rightarrow \Omega^{-2} \{ & R - \xi \{ 6g^{\mu\nu} \nabla_{\mu} \nabla_{\nu} \ln \Omega \\ & - 6(\nabla_{\mu} \ln \Omega)(\nabla^{\mu} \ln \Omega) \} \}. \end{aligned} \quad (20)$$

- On-shell torsion breaks conformal invariance
- Off-shell transformations correspond to different theories of gravity
- Multiple fermions

References

- [E] E. Cartan, "Sur une gnralisation de la notion de courbure de Riemann et les espaces torsion," C. R. Acad. Sci. (Paris) **174**, 593595 (1922).
- [E] T. W. B. Kibble, "Lorentz invariance and the gravitational field," J. Math. Phys. **2**, 212 (1961). doi:10.1063/1.1703702
- [E] D. W. Sciama, "The Physical structure of general relativity," Rev. Mod. Phys. **36**, 463 (1964) Erratum: [Rev. Mod. Phys. **36**, 1103 (1964)].
doi:10.1103/RevModPhys.36.1103
- [E] F. W. Hehl and B. K. Datta, "Nonlinear spinor equation and asymmetric connection in general relativity," J. Math. Phys. **12**, 1334 (1971). doi:10.1063/1.1665738
- [E] Subhasish Chakrabarty, Amitabha Lahiri, Eur. Phys. J. Plus (2018) **133**: 242.

References

-  H. T. Nieh and M. L. Yan, "Quantized Dirac Field in Curved Riemann-cartan Background. 1. Symmetry Properties, Green's Function," *Annals Phys.* **138**, 237 (1982). doi:10.1016/0003-4916(82)90186-5

Thank You!